## The Reduced Density Matrix Method： Current status and open problems

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- Yamashita Makoto
- Michael Overton
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- Ehara Masahiro
- Nakatsuji Hiroshi


## Overview

- Motivation
- Introduction of the RDM method
- Review of results.
- Some open problems (negative results).


## Doing Chemistry from the first principle

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Motivation: Construct a faster, simpler method for Chemistry.

## Candidate: the RDM method

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## The RDM method: 2-RDM as basic variable

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The RDM method: 2-RDM as basic variable

## $\Gamma\left(12 \mid \mathbf{1}^{\prime} \mathbf{2}^{\prime}\right)$

## Basic equation: Schrödinger equation

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Hamiltonian $\boldsymbol{H}$

$$
H=\sum_{j=1}^{N}\left(-\frac{\hbar^{2}}{2 m_{e}} \nabla_{j}^{2}-\sum_{A}^{K} \frac{Z_{A} e^{2}}{4 \pi \epsilon_{0} r_{A j}}\right)+\sum_{i>j} \frac{e^{2}}{4 \pi \epsilon_{0} r_{i j}}
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$$

Schrödinger equation

$$
H \Psi(1,2, \cdots N)=E \Psi(1,2, \cdots N)
$$

Pauli principle

$$
\Psi(\cdots, i, \cdots, j, \cdots)=-\Psi(\cdots, j, \cdots, i, \cdots)
$$

## Solving Schrödinger equation is difficult

We know the basic equation but...

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## We know the basic equation but...



The general theory of quantum mechanics is now almost complete. ... the whole of chemistry are thus completely known, and the difficultly is only that the exact application of these laws leads to equations much too complected to be soluble.

He has never think about computational complexity and it's NP-hard :-)
[Dirac 1929] "Quantum Mechanics of Many-Electron Systems."

## A simpler quantum mechanical method

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A success story: "The Density Functional Theory"
[Hoheberg-Kohn 1964] [Kohn-Sham 1965][Kohn 1998; Nobel Prize]

## Ground state electronic density $\rho(r)$

$\Downarrow$ external potential $\boldsymbol{v}(\boldsymbol{r})$
$\Downarrow$
Hamiltonian $\boldsymbol{H}$
$\Downarrow$
Schrödinger equation
Very difficult functional $\boldsymbol{F}[\rho(r)]$. Practically this is semi-empirical theory.

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> The RDM method!

## The RDM method

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## The second-order reduced density matrix:

[Husimi 1940], [Löwdin 1954], [Mayer 1955], [Coulson 1960], [Nakatsuji 1976]

$$
\Gamma\left(12 \mid 1^{\prime} 2^{\prime}\right)=\binom{N}{2} \int \Psi^{*}(123 \cdots N) \Psi\left(1^{\prime} 2^{\prime} 3 \cdots N\right) d \mu_{3 \cdots N}
$$

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## Comparison

| Method | \# of variable (discritized) | Exact? |
| :---: | :---: | :---: |
| $\boldsymbol{\Psi}$ | $\boldsymbol{N},(\boldsymbol{r}!)$ | Yes |
| $\boldsymbol{\Gamma}\left(\mathbf{1 2 \| \mathbf { 1 } ^ { \prime } \mathbf { 2 } ^ { \prime } )}\right.$ | $\mathbf{4},\left(\boldsymbol{r}^{\mathbf{4}}\right)$ | Yes |

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Do not depend on the size of the system Equivalent to Schrödinger eq. (ground state)

## Definition of 1,2-RDM

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The Hamiltonian contains only 1 and 2-particle interaction.

$$
H=\sum_{i j} v_{j}^{i} a_{i}^{\dagger} a_{j}+\frac{1}{2} \sum_{i_{1} i_{2} j_{1} j_{2}} w_{j_{1} j_{2}}^{i_{1} i_{2}} a_{i_{1}}^{\dagger} a_{i_{2}}^{\dagger} a_{j_{2}} a_{j_{1}}
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$$

The total energy $\boldsymbol{E}$ becomes,

$$
\begin{aligned}
E & =\sum_{i j} v_{j}^{i}\langle\Psi| a_{i}^{\dagger} a_{j}|\Psi\rangle+\frac{1}{2} \sum_{i_{1} i_{j_{1} j_{2}}} w_{j_{1} j_{2}}^{i_{1} i_{2}}\langle\Psi| a_{i_{1}}^{\dagger} a_{i_{2}}^{\dagger} a_{j_{2}} a_{j_{1}}|\Psi\rangle \\
& =\sum_{i j} v_{j}^{i} \gamma_{j}^{i}+\sum_{i_{1} i_{2} j_{1} j_{2}} w_{j_{1} j_{2}}^{i_{1} i_{2}} \Gamma_{j_{1} j_{2}}^{i_{1} i_{2}} .
\end{aligned}
$$

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Here we defined the second-order reduced density matrix $\boldsymbol{\Gamma}_{j_{1} j_{2}}^{i_{1} i_{2}}$ (2-RDM)
$j_{1} j_{2}$

$$
\Gamma_{j_{1} j_{2}}^{i_{1} i_{2}}=\frac{1}{2}\langle\Psi| a_{i_{1}}^{\dagger} a_{i_{2}}^{\dagger} a_{j_{2}} a_{j_{1}}|\Psi\rangle
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$$

and the first-order reduced density matrix $\gamma_{j}^{i}(1-\mathrm{RDM})$

$$
\gamma_{j}^{i}=\langle\Psi| a_{i}^{\dagger} a_{j}|\Psi\rangle
$$

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E_{g} & =\min _{\Psi}\langle\Psi| H|\Psi\rangle \\
& =\min _{\gamma, \Gamma}\left\{\sum_{i j} v_{j}^{i} \gamma_{j}^{i}+\sum_{i_{1} i_{2} j_{1} j_{2}} w_{j_{1} j_{2}}^{i_{i} i_{2}} \Gamma_{j_{1} j_{2}}^{i_{1} i_{2}}\right\}
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$$
\underbrace{N \text {-representability condition }}_{\text {[Coleman 1963] }}
$$

## $N$-representability condition

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$N$-representability condition [Coleman 1963]
$\Gamma\left(\mathbf{1 2} \mid \mathbf{1}^{\prime} \mathbf{2}^{\prime}\right) \rightarrow \Psi(\mathbf{1 2 3} \cdots N)$ or full Density Matrix $\gamma\left(\mathbf{1} \mathbf{1}^{\prime}\right) \rightarrow \Psi(\mathbf{1 2 3} \cdots N)$ or full Density Matrix

## Complete $N$-representability condition is very difficult

Computational complexity for complete $N$-representability

- Basically every $N$-rep. condition is linear constraint.
- Optimization over diagonal $N$-representable 2-RDM: NP-hard [Deza, Laurent 1997]
- Decision problem "is this 2-RDM $N$-representable?" QMA-complete, NP-hard [Liu 2007]
- Coleman's algorithm [Beste, et al 2002] [Coleman 2000].

Very difficult and hopeless :-(
[See also "Complexity and Electronic Structure Theory" B.J. Braams 741 PT11 16:40 Monday]

## Approximate $N$-representability condition

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- T1, T2, T2', ( $\overline{\boldsymbol{T}} \mathbf{2}$ )-condition [Zhao et al. 2004], [Erdahl 1978] [Braams et al 2007] [Mazziotti 2006, 2007]


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- Davidson's inequality [Davidson 1969][Ayers et al. 2006]
- Construction of 2-particle density [Pistol 2004, 2006]


## PSD type $N$-representability conditions

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$\boldsymbol{P}, \boldsymbol{Q}, \boldsymbol{G}, \boldsymbol{T 1}, \boldsymbol{T 2}$-matrix are all positive semidefinite $\leftrightarrow$ eigenvalues $\lambda_{i} \geq \mathbf{0}$

$$
U^{\dagger} \Gamma U=\left[\begin{array}{llll}
\lambda_{1} & & & 0 \\
& \lambda_{2} & & \\
& & \ddots & \\
0 & & & \lambda_{n}
\end{array}\right] \geq 0
$$

First application to Be atom
[Garrod et al 1975, 1976]
Calculation methods are not very well studied...

## Realization of the RDM method for atoms and molecules

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$$
\begin{gathered}
\boldsymbol{E}_{\mathbf{g}}=\underset{\boldsymbol{\Gamma} \in \mathcal{P}}{\operatorname{Min}} \operatorname{Tr} \boldsymbol{H} \boldsymbol{\Gamma} \\
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[Nakata-Nakatsuji-Ehara-Fukuda-Nakata-Fujisawa 2001]
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## Semidifinite programming

We solved exactly for the first time!

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## The ground state energy of atoms and molecules [Nakata et al 2008]

| System | State | $\boldsymbol{N} \quad \boldsymbol{r}$ | $\Delta \mathrm{E}_{\text {GT1T2 }}$ | $\Delta \mathrm{E}_{\text {GT1T2 }}$ | $\Delta \mathrm{E}_{\text {CCSD }}(T)$ | $\Delta \mathrm{E}_{\text {HF }}$ | $\mathrm{E}_{F C I}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| C | ${ }^{3} \mathrm{P}$ | 620 | -0.0004 | -0.0001 | +0.00016 | +0.05202 | -37.73653 |
| O | ${ }^{1} \mathrm{D}$ | 820 | -0.0013 | -0.0012 | +0.00279 | +0.10878 | -74.78733 |
| Ne | ${ }^{1} S$ | 1020 | -0.0002 | -0.0001 | -0.00005 | +0.11645 | -128.63881 |
| O | ${ }^{2} \Pi_{g}$ | 1520 | -0.0022 | -0.0020 | +0.00325 | +0.17074 | -148.79339 |
| BH | ${ }^{1} \Sigma^{+}$ | 624 | -0.0001 | -0.0001 | +0.00030 | +0.07398 | -25.18766 |
| CH | ${ }^{2} \Pi_{r}$ | 724 | -0.0008 | -0.0003 | +0.00031 | +0.07895 | -38.33735 |
| NH | ${ }^{1} \Delta$ | 824 | -0.0005 | -0.0004 | +0.00437 | +0.11495 | -54.96440 |
| HF | ${ }^{1} \Sigma^{+}$ | 1424 | -0.0003 | -0.0003 | +0.00032 | +0.13834 | -100.16031 |
| $\mathrm{SiH}_{4}$ | ${ }^{1} A_{1}$ | 1826 | -0.0002 | -0.0002 | +0.00018 | +0.07311 | -290.28490 |
| $\mathrm{F}^{-}$ | ${ }^{1} S$ | 1026 | -0.0003 | -0.0003 | +0.00067 | +0.15427 | -99.59712 |
| P | ${ }^{4} S$ | 1526 | -0.0001 | -0.0000 | +0.00003 | +0.01908 | -340.70802 |
| $\mathrm{H}_{2} \mathrm{O}$ | ${ }^{1} A_{1}$ | 1028 | -0.0004 | -0.0004 | +0.00055 | +0.14645 | -76.15576 |

$\boldsymbol{G T 1 T 2}$ : $\quad$ The RDM method ( $\boldsymbol{P}, \mathbf{Q}, \boldsymbol{G}, \boldsymbol{T 1}$ and $\boldsymbol{T 2}$ conditions)
GT1T2' : The RDM method ( $P, Q, G, T 1$ and $T 2^{\prime}$ conditions)
$\operatorname{CCSD}(\mathrm{T})$ : Coupled cluster singles and doubles with perturbation treatment of triples
HF : Hartree-Fock
FCl : FullCl

## Application to potential energy curve

- Dissociation curve of $\mathbf{N}_{2}$ (triple bond) the world first result. [Nakata-Nakatsuji-Ehara 2002]



## Summary of typical results

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## Typical results

$N$-rep. Correlation energy(\%) dissociation limit
$P Q G \quad 100$ ~ 130\%
PQGT1T2' $\operatorname{CCSD}(\mathrm{T})$
$100 \sim 101 \%$ $100 \sim 101 \%$
([Nakata-Nakatuji-Ehara-Fukuda-Nakata-Fujisawa 2001] [Nakata et al.2002][Zhao et al. 2004][Mazziotti et al. 2002,3,4,5,6,7,8,9][Eric et al 2006][Fukuda et al 2007][Nakata et al 2008] )
[See also "Variational Detremination of the Two Particle Reduced Density Matrix", B. Verstichel 1319 PTP Tuesday]

## Essentials of [Nakata et al. 2001, 2002] with $P, Q, G$-condition

## JCP 114,8282(2001), JCP 116, 5432(2002).

- Exact minimization was performed for the first time. Since mid of 1990, solvers are developed mathematical programming community. Among them the SDPA the fastest, and exact. Everyone is doing SDP without aware of it :)
- G-condition is mandatory for chemistry: related to AGP wf (corr. ene. 100\%~130\%)
- $P, Q$ conditions are not sufficient (corr. ene. $200 \sim \mathbf{8 0 0 \%}$ )
- The 3rd, 4th order approx. are not essential
- In 2000 Erdahl and Jin:"The poor quality of the second-order estimates is consistent with the work of Mihailovic and Rosina..."
- In 2001, Mazziotti and Erdahl: Lipkin model. Table IV: correlation energies are 478-5238\% with $\boldsymbol{P}, \boldsymbol{Q}, \boldsymbol{G}$ conditions (in 2005 silently corrected by Hammond, Mazziotti and Erdahl(?)).


## Importance of [Zhao-Braams-Fukuda-Percus-Overton 2004] with

## T1, T2-conditions

JCP 120, 2095(2004).

- Formulation of $\boldsymbol{T 1}$ and $\boldsymbol{T 2}$ conditions
- Total energies became comparable to CCSD(T)
- We can do Chemistry with these conditions!!!
- Not a trivial formulation.
- Dual SDP formulation

Lager calculation is possible

- They gave us a hope to realization of the RDM method


## How large we can calculate until now?

"size" is how many active orbitals are used

| Who | Year | Size | $N$-rep. | system |
| :---: | :---: | :---: | :---: | :---: |
| Garrod et al. | 1976 | ? | PQG | Be |
| Erdahl | 1979 | ? | $P Q G$ | $\mathrm{He}_{2}$ |
| Nakata et al. | 2001 | 8 | PQG | $\mathrm{H}_{3} \mathrm{O}^{+}$ |
| Nakata et al. | 2002 | 8 | $P Q G$ | $\mathrm{C}_{2}, \mathrm{CO}, \mathrm{N}_{2}$ |
| Zhao et. al | 2004 | 10 | PQGT1T2 | $\mathrm{H}_{3} \mathrm{O}^{+}$ |
| Mazziotti | 2004 | 14-18 | $P Q G$ | $\mathrm{N}_{2}, \mathrm{H}_{6}$ |
| Gidofalvi et.al. | 2005 | 14-20(?) | $P Q G$ | $\mathrm{N}_{2}, \mathrm{C}_{2}$ |
| Eric et. al | 2006 | 15 | $P Q G$ | $\mathrm{NH}_{3}$ |
| Hammond et. al | 2006 | 14 | PQGT2 | Hubbard mod |
| Fukuda et al. | 2007 | 13 | PQGT1T2 | $\mathrm{CH}_{3} \mathrm{~N}$ |
| Nakata et al. | 2008 | 14 | PQGT1T2' | $\mathrm{H}_{2} \mathrm{O}$ |
| Greenman et al. | 2009 | 8 | PQGT2 | CAS(Benzene |
| Greenman et al. | 2009 | 12-24 | $P Q G$ | CAS(Pentacyn |

## Open problems

- How many iterations are needed?
- Size consistency or extensively.
- Degeneracy.
- Diagonal representability and excited states.


## How many iterations are needed?

How many iterations are required by

- primal-dual interior-point method (PDIPM) or
- Monteiro-Bruner method (RRSDP) [Mazziotti 2004]

|  | $\boldsymbol{P}, \boldsymbol{Q}$, and $\boldsymbol{G}$ |  |  | $\boldsymbol{P}, \boldsymbol{Q}, \boldsymbol{G}, \boldsymbol{T} \mathbf{1}, \boldsymbol{T} \mathbf{2}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| algorithm | flops | \# iterations | memory | flops | \# iterations | memory |
| PDIPM | $\boldsymbol{r}^{12}$ | $\boldsymbol{r} \ln \boldsymbol{\varepsilon}^{-1}$ | $\boldsymbol{r}^{\mathbf{8}}$ | $\boldsymbol{r}^{12}$ | $\boldsymbol{r}^{\mathbf{3 / 2}} \ln \boldsymbol{\varepsilon}^{-1}$ | $\boldsymbol{r}^{\mathbf{8}}$ |
| RRSDP | $\boldsymbol{r}^{\mathbf{6}}$ | none | $\boldsymbol{r}^{4}$ | $\boldsymbol{r}^{9}$ | none | $\boldsymbol{r}^{\mathbf{6}}$ |

Note: when we stop the iteration is a big problem

## Lacking size consistency and/or extensively

- $\boldsymbol{G}$-condition servers as a necessary condition to size consistent and/or extensively [Nakata et al. 2002] $\lim _{\left|1-1^{\prime}\right| \rightarrow \infty} F\left(1 \mid 1^{\prime}\right)=\left\langle n(1) n\left(1^{\prime}\right)\right\rangle-\langle n(1)\rangle\left\langle n\left(1^{\prime}\right)\right\rangle \geq 0 \quad$ this should be zero
- PQG are not consistent: $\mathbf{C}_{\mathbf{2}}, \mathbf{C O} 33.3 \mathrm{mH}$, and 5.8 mH .
- PQG seems to be size-extensive [Mazziotti et al. 2005].
- $\mathbf{H}_{2} \mathbf{O}, \mathbf{N H}_{\mathbf{3}}$ and $\mathbf{N}_{\mathbf{2}}$ are size consistent. (if molecules dissociates to $\mathbf{H}$, it seems to be size consistent) [JCP 2002]
- Reduce spin number by adding ghost hydrogen atoms [Hammond Mazziotti 2005].
- See also "Dissociation Curves from ...Density Matrices:..." P.Bultink et al. 1552 PT3 [Bultink et al. PCCP 20091


## Ensemble $N$-representability problem: zero dipole

$\mathbf{H}_{3}$ is a meta stable molecule and have two dipole moments at the ground state.


- SDP solver calculates ensemble average of these two dipole moments.
- Adding small perturbation can remove this.
[Fukuda et al 2007][Nakata et al. 2008]


## non- $N$-representability: Spin degeneracy

- Spin degeneracy will be observed when $\boldsymbol{S} \neq \mathbf{0}$.
- E.g., $\boldsymbol{S}=\mathbf{1}$, then the total energies of three states $S_{z}=1,0,-1$ will degenerate.
- Including T1, T2' recovers, but not complete.
- This is a source of non-size consistency or size-extensively.

| System | State | basis | $\operatorname{corr}^{\text {PQG }}$ | corr ${ }_{\text {T1 }}$ | $\operatorname{corr} E_{T 12^{\prime}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| C | ${ }^{3} P_{1}$ | double- $\zeta$ | 107.5\% | 105.9\% | 100.1\% |
| C | ${ }^{3} P_{0}$ | double- $\zeta$ | 133.4\% | 126.0\% | 103.9\% |
| 0 | ${ }^{3} P_{1}$ | double- $\zeta$ | 117.4\% | 109.4\% | 101.1\% |
| 0 | ${ }^{3} P_{0}$ | double- $\zeta$ | 134.4\% | 127.4\% | 102.3\% |

corrE means correlation energy error in percentage
[Nakata et al. 2008]

## Properties becomes better like total energy?

- The total energy becomes higher when we add a new $N$-representability conditions:

$$
E_{P Q} \leq E_{P Q G} \leq E_{P Q G T 1} \leq E_{P Q G T 1 T 2} \leq \cdots \leq E_{\mathrm{FCI}}
$$

No such kind of conditions for properties, but we USUALLY have a following sequence: [Nakata et al 2008]

$$
\left|\mu_{P Q G}-\mu_{\mathrm{FCI}}\right| \leq\left|\mu_{P Q G T 1}-\mu_{\mathrm{FCI}}\right| \leq\left|\mu_{P Q G T 2}-\mu_{\mathrm{FCI}}\right| \leq 0
$$

$\mu$ is the dipole moment of atoms and molecules. Can be arbitrary operators up to 2-body.

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- GMP (GNU multiple precision) $\Rightarrow$ necessity of highly accurate solver, using multiple precision arithmetic (SDPA-GMP) http://sdpa.indsys.chuo-u.ac.jp/sdpa/ GNU Public License


## SDPA-GMP and Hubbard model

## The 1D Hubbard model with high correlation limit $|\boldsymbol{U} / \boldsymbol{t}| \rightarrow \infty$ : All states are almost degenerated.

The ground state energies of 1D Hubbard model
PBC, \# of sites:4, \# of electrons: 4, spin 0

| U/t | SDPA (16 digits) | SDPA-GMP (60 digits) | fullCl |
| :---: | :---: | :---: | :---: |
| 10000.0 | 0 | -1.1999998800000251 $\times 10^{-3}$ | -1.199999880 $\times 10^{-3}$ |
| 1000.0 | -1.2 $\times 10^{-2}$ | -1.1999880002507934 $\times 10^{-2}$ | $-1.1999880002 \times 10^{-2}$ |
| 100.0 | $-1.1991 \times 10^{-1}$ | -1.1988025013717993 $\times 10^{-1}$ | -1.19880248946 $\times 10^{-1}$ |
| 10.0 | -1.1000 | -1.0999400441222934 | -1.099877772750 |
| 1.0 | -3.3417 | -3.3416748070259956 | -3.340847617248 |
| PBC, \# of sites:6, \# of electrons: 6, spin 0 |  |  |  |
| U/t | SDPA (16 digits) | SDPA-GMP (60 digits) | fullCl |
| 10000.0 | 0 | -1.7249951195749525 $\times 10^{-3}$ | -1.721110121 $\times 10^{-3}$ |
| 1000.0 | $-1 \times 10^{-2}$ | $-1.7255360310431304 \times 10^{-2}$ | $-1.7211034713 \times 10^{-2}$ |
| 100.0 | $-1.730 \times 10^{-1}$ | $-1.7302157140594339 \times 10^{-1}$ | -1.72043338097 $\times 10^{-1}$ |
| 10.0 | -1.6954 | -1.6953843276854447 | -1.664362733287 |
| 1.0 | -6.6012 | -6.6012042217806286 | -6.601158293375 |

## Excited states systems

- RPA: [Garrod et al 1980] [Mazziotti 2003] $\rightarrow$ can go beyond?
- McDonald's variational principle [Nakata et al 2006][Erdahl Grudziński 1978][Yasuda 2002] $\rightarrow$ Excited states of one-particle Hamiltonian: essentially difficult problem. Related to 2-particle density, too.


## Open problems

- Urgent How to recover size extensively and/or consistency
- Urgent Large scale semi-definite programming solver.
- Degeneracy problems.
- Understanding $N$-representability conditions: Physical and Chemical meaning.
- How to find "a good" $\boldsymbol{N}$-representability conditions.


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> WELCOME TO THE JUNGLE!

## Missing topics

I missed many very very interesting areas

- Density equation / contracted Schrödinger equation and its variants.
- 1-RDM functional theories.
- 2-RDM parametrization theories.
- 2-particle density functional theories.


## How large these SDP are?

## \# of constraints

$r$ constraints block

| 24 | 15018 | $2520 \times 2,792 \times 4,288 \times 1,220 \times 2$ |
| :---: | :---: | :---: |
| 26 | 20709 | $3211 \times 2,1014 \times 4,338 \times 1,286 \times 2$ |

Elapsed time using Itanium 2 ( 1.3 GHz ) 1 node 4 processors. System, State, Basis $N$-rep. $\quad r \quad$ Time $\#$ of nodes $\mathbf{S i H}_{4},{ }^{1} \boldsymbol{A}_{\mathbf{1}}$, STO-6G PQGT1T2 $26 \quad 5.1$ days $\quad 16$ $\mathbf{H}_{\mathbf{2}} \mathbf{O},{ }^{1} \boldsymbol{A}_{\mathbf{1}}$, double- $\boldsymbol{\zeta} \quad \boldsymbol{P Q G} \quad 28 \quad 2.2$ hours $\quad 8$ $\mathbf{H}_{\mathbf{2}} \mathbf{O},{ }^{1} \boldsymbol{A}_{1}$, double- $\zeta \quad$ PQGT1T2 $28 \quad 20$ days $\quad 8$ $\begin{array}{lllll}\mathbf{H}_{\mathbf{2}} \mathbf{O},{ }^{1} \boldsymbol{A}_{\mathbf{1}}, \text { double- } \zeta & \boldsymbol{P Q G T 1 T 2} & 28 & 24 \text { days } & 8\end{array}$

